

# GQuEST

## Introduction and Overview

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# GQuEST collaboration at Caltech September 2023



# A Brief History of GQuEST

## Gravity from the Quantum Entanglement of Space-Time

Kathryn et al. describe quantum gravity observables in interferometers. Rana and Lee's experience at LIGO and the Holometer motivate the design.

Fermilab has Holometer equipment, experience with: cryogenics, optics, digital signal processing and projects.

The GQuEST collaboration has funding from

- Department of Energy: Office of Science
- Heising-Simons Foundation

to design, build, and operate the experiment. We received one year of funding from DOE with more expected.

Current focus is to build a TECHNOLOGY DEMONSTRATOR

5 year time frame



# “Elevator Speech”

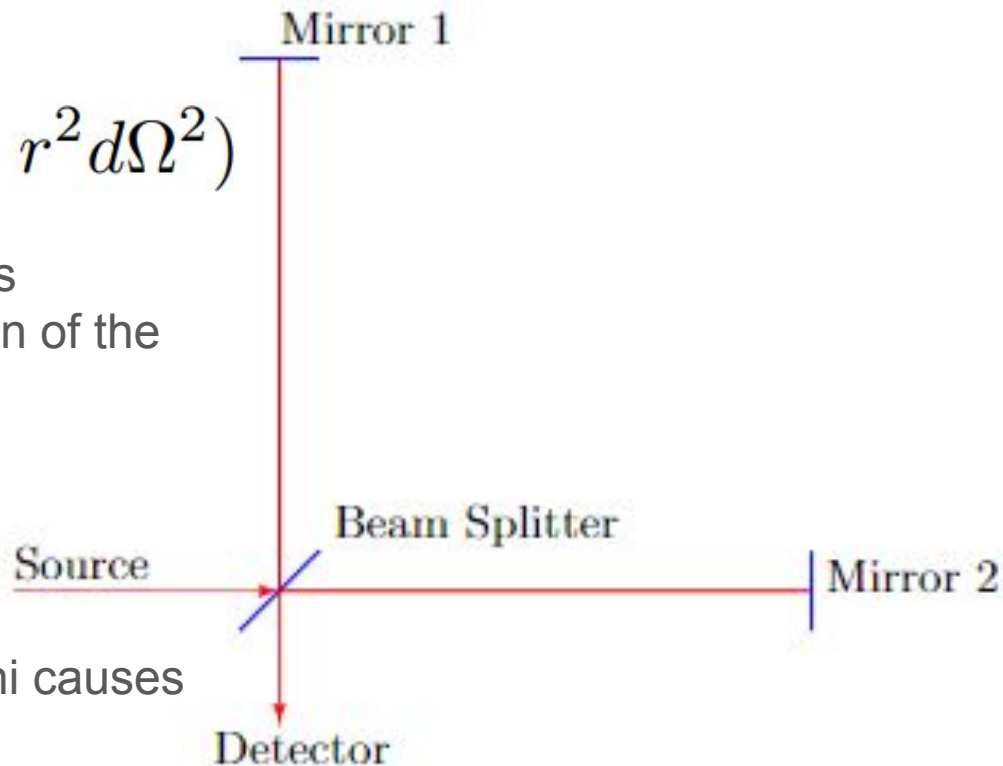
$$ds^2 = -dt^2 + (1 - \phi)(dr^2 + r^2 d\Omega^2)$$

IFOs are operated by setting the distances to the two mirrors such that a small fraction of the source power reaches the detector.

Changing the phase difference changes the power on the detector

In this version of the metric, a non-zero  $\phi$  causes phase change that we detect.

This detection is at the input light's reference frequency  $c/\lambda$  and also in sidebands.



**We focus on the sidebands!**

# Fiducial GQuEST Design

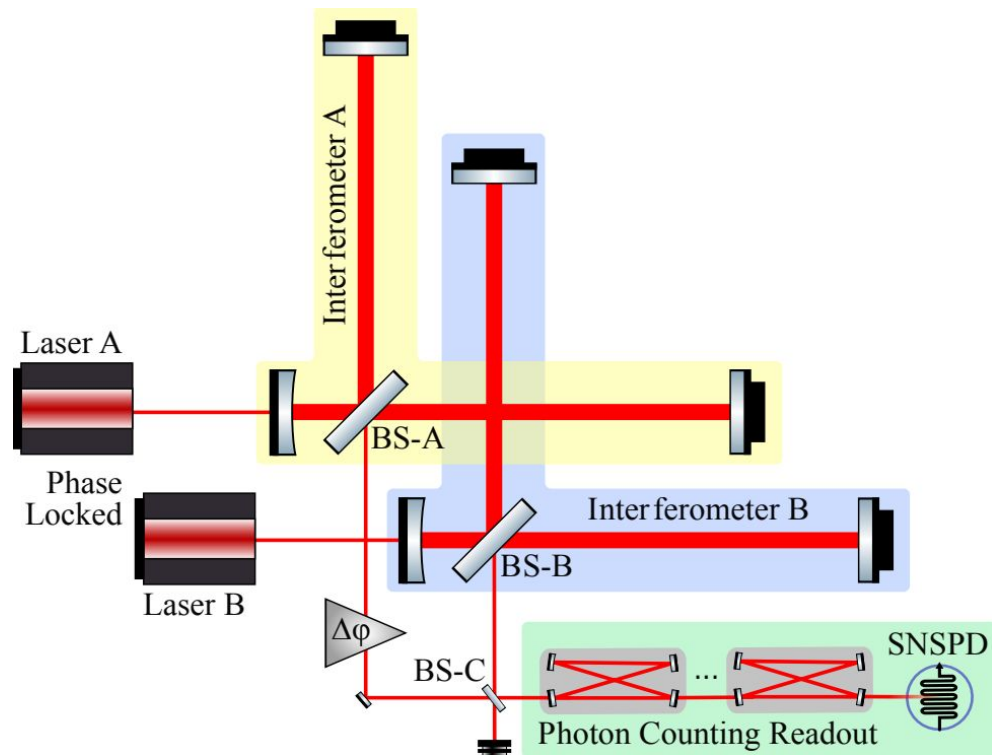
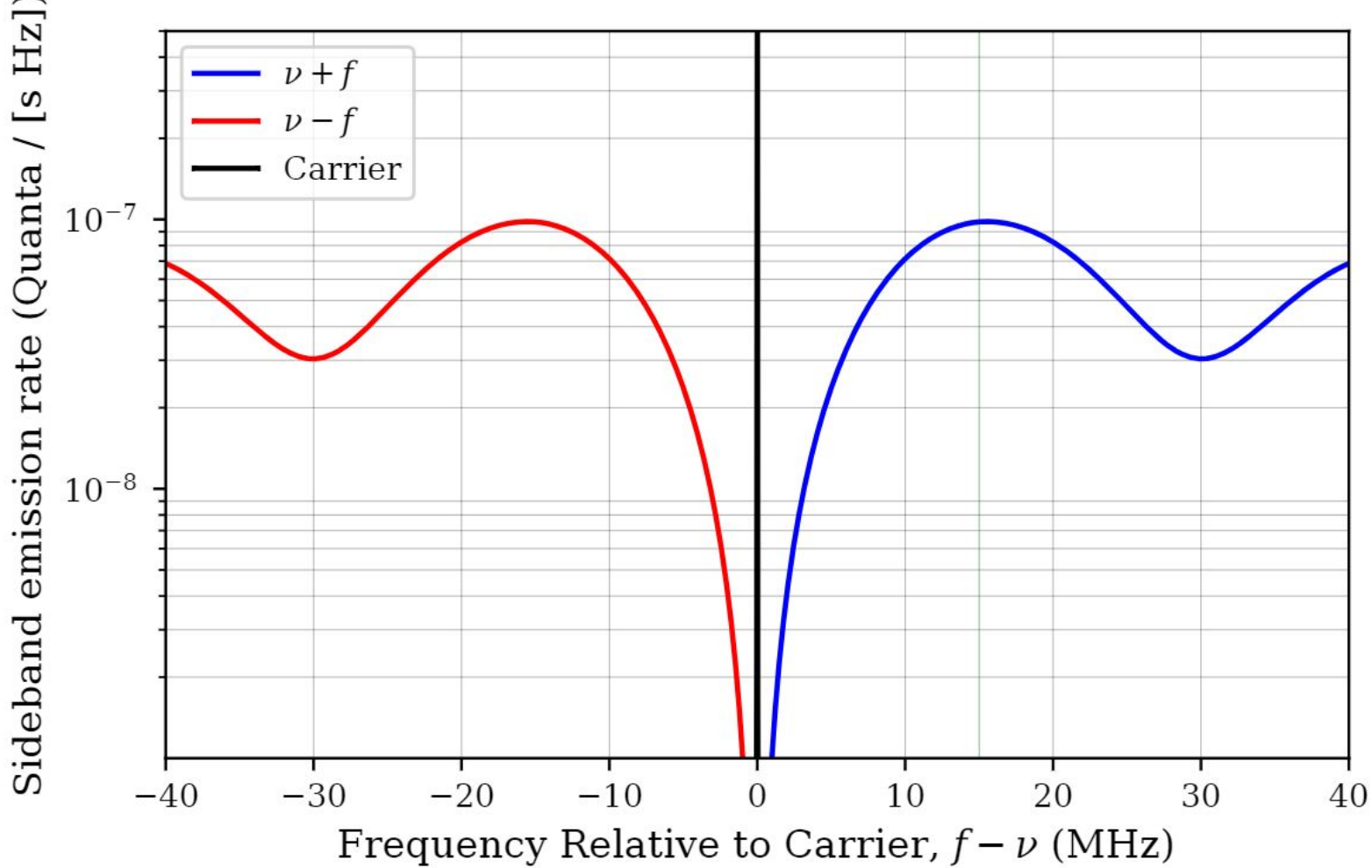


TABLE I. Parameters of the fiducial GQuEST design. The noise spectral densities are evaluated at 17.6 MHz.

parameter	symbol	value
Geotropic fluct. scale parameter	$\alpha$	$\mathcal{O}(1)$
IFO arm length	$L$	5 m
Power on beamsplitter	$P_{BS}$	10 kW
Laser wavelength	$\lambda$	1550 nm
Laser frequency	$\nu$	193.4 THz
Nominal filter offset frequency	$\epsilon_c$	17.6 MHz
Filter bandwidth (FWHM)	$\Delta\epsilon$	25 kHz
Twin IFO separation	$L_s$	1.5 m
IFO inter-arm angle	$\Theta$	$90^\circ$
Signal Spectral Density (peak)	$\bar{S}_L^\phi$	$(3 \cdot 10^{-22} \text{ m}/\sqrt{\text{Hz}})^2$
Thermal Noise Spectral Density	$\bar{S}_L^c$	$(10^{-21} \text{ m}/\sqrt{\text{Hz}})^2$
Shot Noise Spectral Density	$\bar{S}_L^q$	$(6 \cdot 10^{-19} \text{ m}/\sqrt{\text{Hz}})^2$
Filtered signal photon flux	$\dot{N}_{\text{pass}}^\phi$	$1.4 \cdot 10^{-3} \text{ Hz}$
Filtered classical noise photon flux	$\dot{N}_{\text{pass}}^c$	$1.6 \cdot 10^{-2} \text{ Hz}$
Photon Detector Dark Count Rate	$\dot{N}^d$	$< 10^{-3} \text{ Hz}$
Observation time for $5\sigma$ test for $\alpha = 1$	$T$	$\mathcal{O}(10^5) \text{ s}$



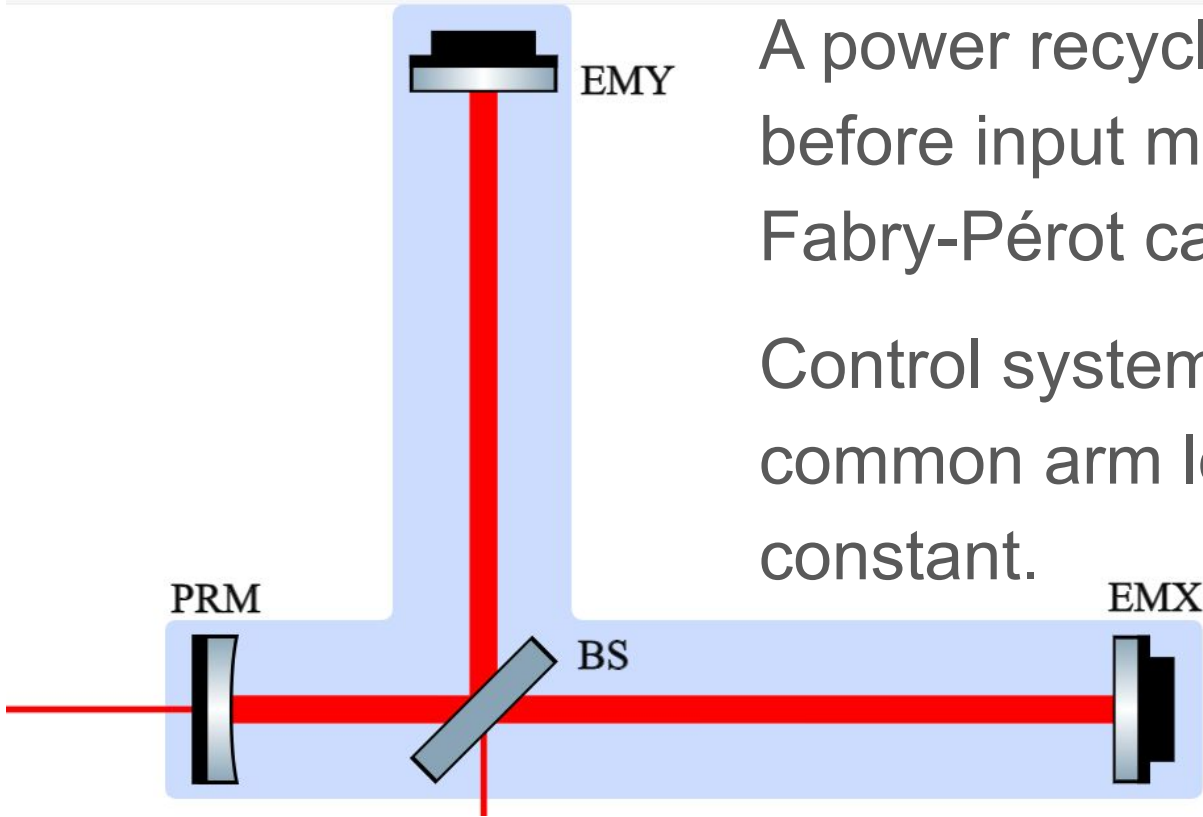


## Experimental Considerations: Physical Stability

- On paper the length difference (and hence phase difference) is constant.
- In practice, everything moves! Mechanical vibration, laser frequency, effects in mirror surface and substrate, ....
- We keep the differential arm length (aka DARM) constant with feedback systems to keep the output power at a constant level.



# Experimental Considerations: More photons on target



A power recycling mirror (PRM) before input makes the IFO a Fabry-Pérot cavity.

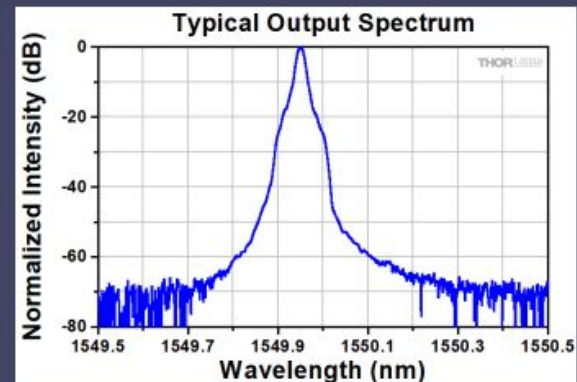
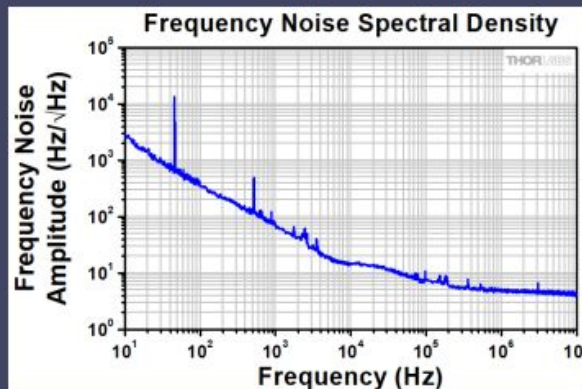
Control system keeps the common arm length (CARM) constant.



# Experimental Considerations: laser phase noise $\rightarrow$ sidebands

- Input laser phase noise (White):  $10^{-7}$  rad/ $\sqrt{\text{Hz}}$
- Laser Power: 10 W
- Shot noise with cavity power:  $6.2 \cdot 10^{-19}$  m/ $\sqrt{\text{Hz}}$

$\Delta$  is 0.1 nm  $\sim$  12 MHz



See Ian MacMillan's talk on Laser Noise

## Experimental Considerations: homodyne readout

The “standard quantum limit” arises from poisson noise on the number of photons/second on the beam splitter.

Expressed as (effective) length changes:

$$\overline{S}_L^q = \frac{\hbar c}{2kP_{\text{BS}}} \approx \left( 6.2 \cdot 10^{-19} \frac{\text{m}}{\sqrt{\text{Hz}}} \right)^2 \left( \frac{10 \text{ kW}}{P_{\text{BS}}} \right) \left( \frac{\lambda}{1550 \text{ nm}} \right)$$

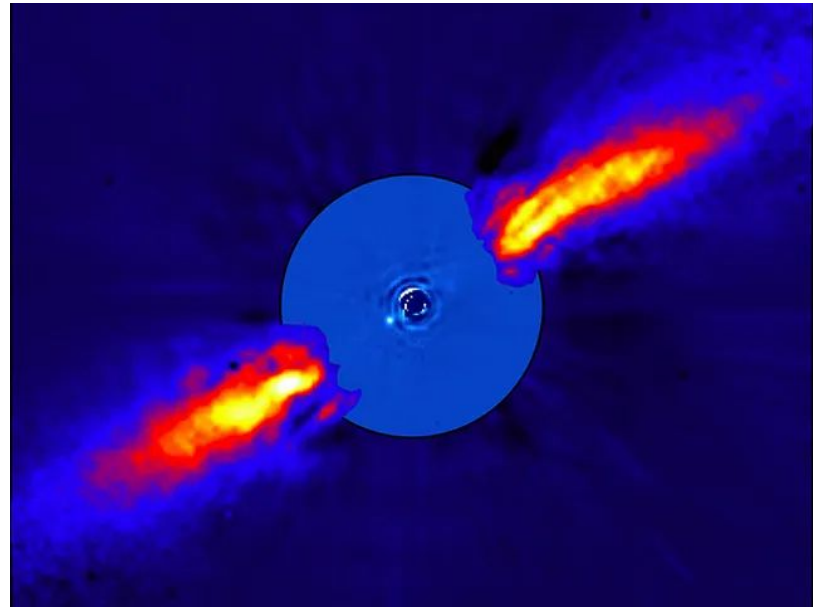
The expected signal ( $3 \cdot 10^{-22} \text{ m}/\text{rtHz}$ ) is  $\ll$  this noise.

See Sander Vermeulen's talk on Interferometry

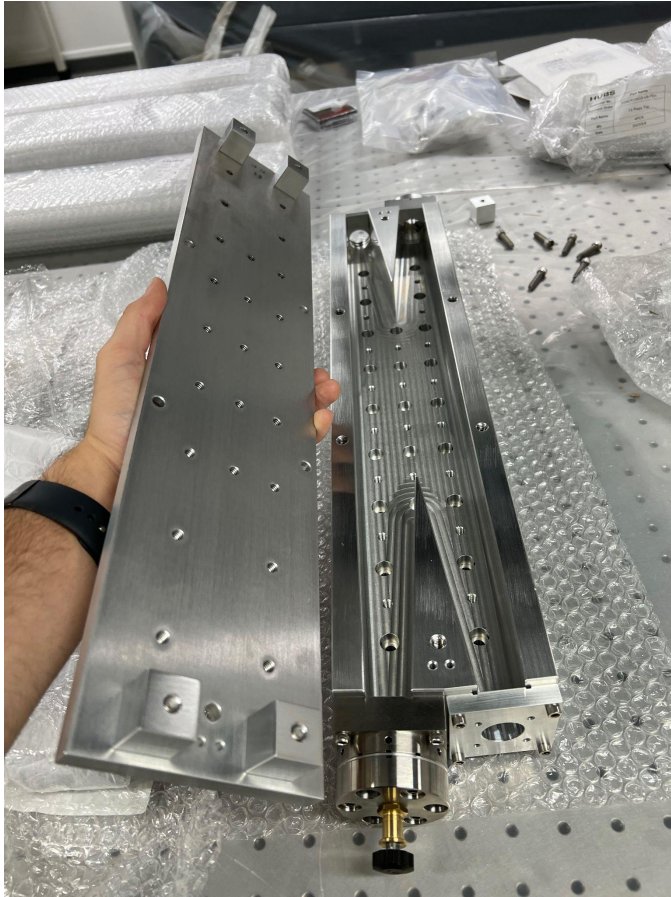
# Experimental Considerations

S/N in the SIDEBAND scales more favorably, but the ratio of power (near  $\lambda$ )/(sidebands) is very large.

An analogy in the spatial domain is coronagraph imaging of exoplanetary systems.



# Experimental Considerations: Output Filters



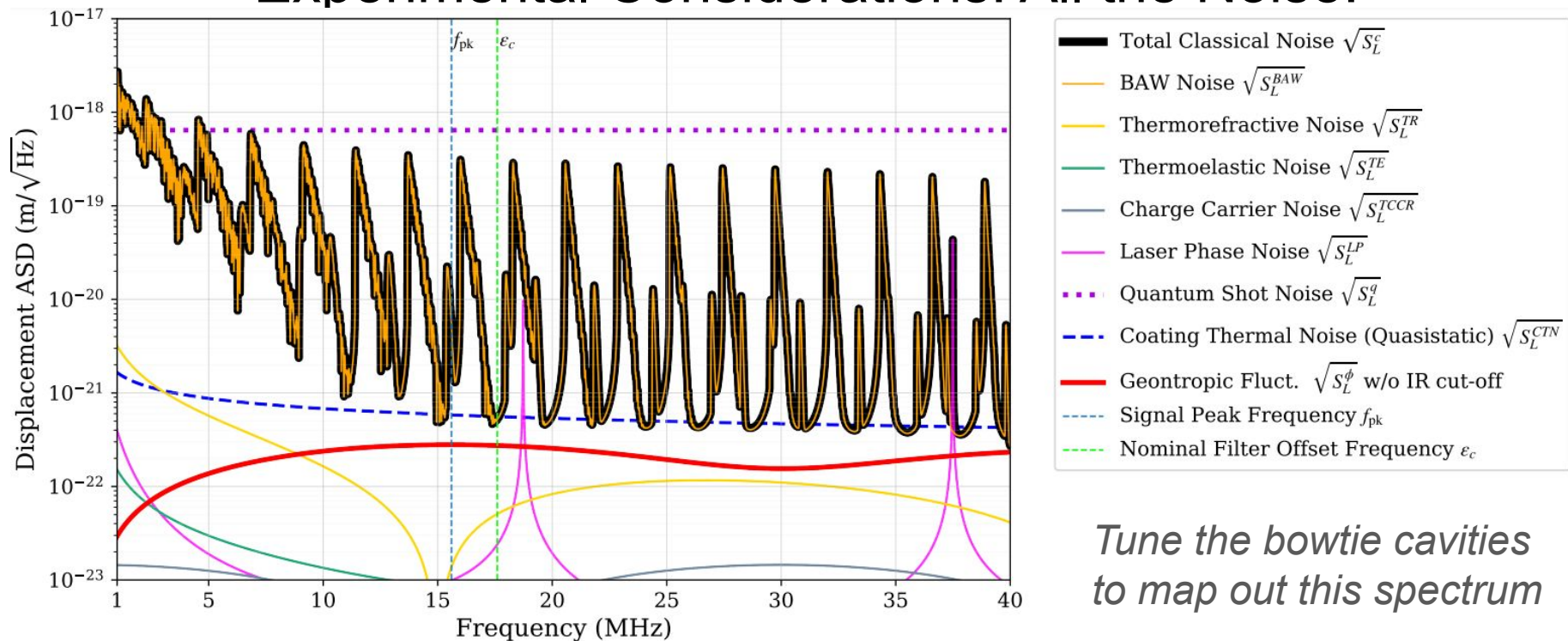
Non-signal light suppressed by  $\sim 60$  dB  
(factor of  $10^6$ ) in the frequency domain

The “bowtie cavities” resonance  
frequency can be actively tuned to be  
from 8 to 40 MHz

Four cavities in series yield 240 dB  
(factor of  $10^{24}$ ) suppression

See Torrey Cullen’s talk on Technology Demonstrator

# Experimental Considerations: All the Noise!



*Tune the bowtie cavities  
to map out this spectrum*

Mirrors are not perfect: the substrate is a drum; coatings have thermal noise

See Daniel Grass' talk on Thermal Noise



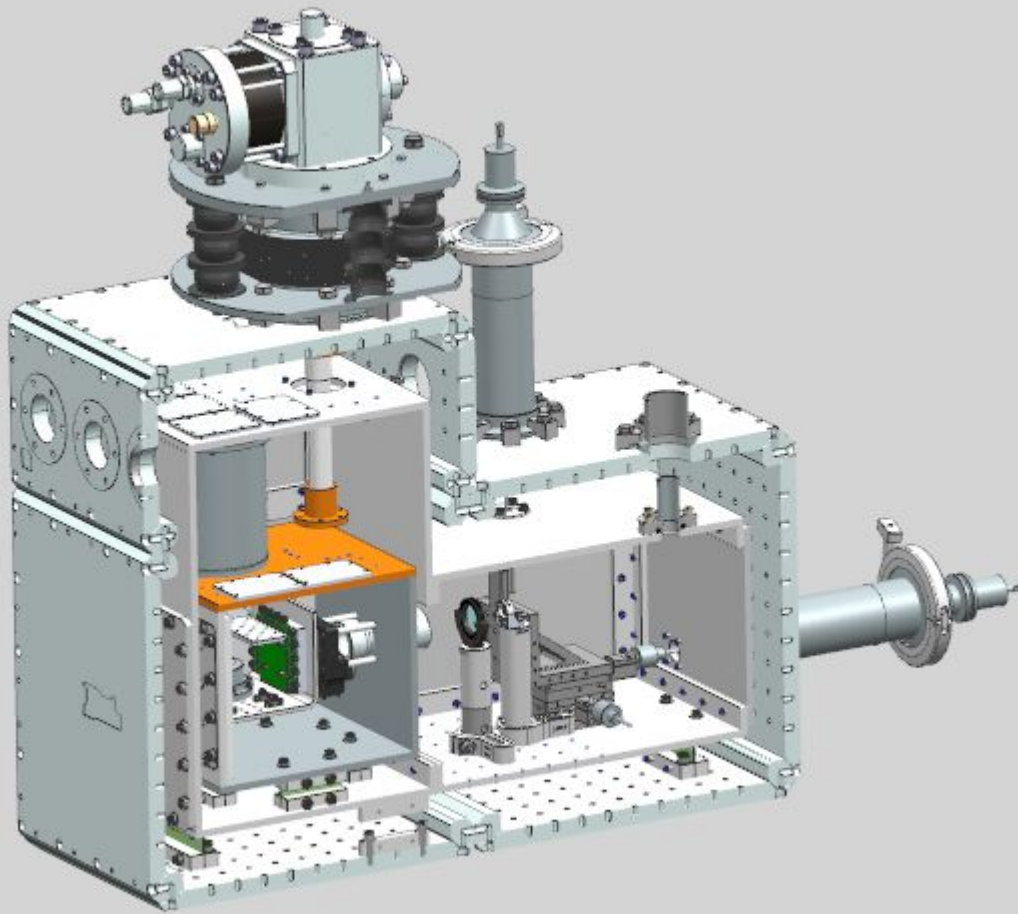
# Experimental Considerations: Photon Counter

- JPL is providing SNSPD counters: Superconducting Nanowire Single Photon Detector.
- Thermal Radiation photons would overwhelm the rate. Sources are:
  - Near the detector (so place in cold dewar)
  - Shining from IFO optics (so couple with fiber and optical filters)

See Alex Ramirez's talk on Photon Counting







Dewar under construction at Fermilab to hold:

- SNSPD detector,
- optical filters
- thermal shields
- alignment

# Science Operations

Map out spectrum by repeating these steps:

- Choose a frequency
- Count photons for  $\sim 1$  hour to 24 hours.

Characterize the EXCESS:

- Modulate signal by changing  $\Delta\varphi$ 
  - This modulation can be time-dependent, blinded
- Change  $L$ 
  - Technology Demonstrator will have a small  $L$



Details on Friday:

Sander Vermeulen - Interferometry Basics and current limits

Ian MacMillan - Laser noise

Daniel Grass - Thermal noise

Torrey Cullen - Technology demonstrator and filter cavities

Alex Ramirez - Detecting Photons One At A Time

Lee McCuller - Future quantum experiment possibilities

